Feynman path integrals for exponents of polynomially growing functionals

E.T. SHAVGULIDZE*)

Department of Mechanics and Mathematics, Lomonosov Moscow State University, Vorobievy Gory, 119899 Moscow, Russia

A general class of functional integrals of the exponents of polynomially growing functionals on the Hilbert space be studied. A representation formula by integrals for Gaussian measures is given for this class of functional integrals. These results are applied to provide a rigorous Feynman path integral representations for the solutions of the time dependent Schrodinger equations with a polynomially growing potentials (it is possible with alternating signs). Special self-adjoint extensions for Schrodinger differential operators with a polynomially growing potentials are obtained.

Let *H* be a real separable Hilbert space, with inner product (\cdot, \cdot) , and norm $\|\cdot\|$, *T* be a self-adjoint strictly positive trace class operator.

Let $e_1, e_2, \ldots, e_n, \ldots$ be orthonormal base of Hilbert space H such that $Te_n = \lambda_n e_n, \lambda_n > 0, \sum_{n=1}^{\infty} \lambda_n < +\infty$.

Denoted by H_n the span of the first *n* vectors e_1, e_2, \ldots, e_n . Let P_n be the orthogonal projector from the Hilbert space *H* onto the subspace H_n .

Denoted by μ the Gauss measure on the Hilbert space H with correlation operator T and the zero mean value. The Gauss measure μ is countably additive.

Let $Q_{2l}: H \times \cdots \times H \to R$ be 2*l*-linear symmetric continuous function such that $Q_{2l}(x, \ldots, x) > 0$ for all $x \in H, x \neq 0$, and $p(x) = (Q_{2l}(x, \ldots, x))^{1/(2l)}$.

Definition. Let $v \in \mathbf{C}$, $v \neq 0$, $\operatorname{Re} v \geq 0$ be a constant and $f : H \to \mathbf{C}$ be a continuous function such that for all natural number n it exists Fresnel integral

$$\int_{H_n} f(x) \exp\left(-\frac{\upsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x = \lim_{\epsilon \to +0} \int_{H_n} f(x) \exp\left(-\frac{\upsilon + \epsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x.$$

We define the Feynman integral

$$\int_{H} f(x) \exp\left(-\frac{\upsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x$$

as the limit

$$I = \lim_{n \to \infty} \frac{\int_{H_n} f(x) e^{-v(T^{-1}x,x)/2} dx}{\int_{H_n} e^{-v(T^{-1}x,x)/2} dx}$$

 $\begin{array}{l} \text{We put } \int_{H}f(x)\exp\left(-\frac{1}{2}\,\upsilon(T^{-1}x,x)\right)\mathrm{d}x=I.\\ \text{It is easy to see that} \end{array}$

$$\int_{H_n} \exp\left(-\frac{\upsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x = \sqrt{\left(\frac{2\pi}{\upsilon}\right)^n \lambda_1 \dots \lambda_n}$$

^{*)} E-mail:SHAV@MECH.MATH.MSU.SU

and

$$\int_{H} f(x) \exp\left(-\frac{\upsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x = \lim_{n \to \infty} \frac{\int_{H_n} f(x) \exp\left(-\frac{\upsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x}{\sqrt{\left(\frac{2\pi}{\upsilon}\right)^n \lambda_1 \dots \lambda_n}}$$

•

Let F(H) be the space of all continuous functions $f: H \to \mathbb{C}$ such that for all $h_1, h_2 \in H$ the function $\varphi(\alpha, \beta) = f(\alpha h_1 + \beta h_2)$ have analytic continuation $\mathbb{C} \times \mathbb{C}$ and for all $\alpha, \beta \in \mathbb{C}$. It is true

$$|\varphi(\alpha,\beta)| \le c_1 \exp\left(c_2 \left[|\alpha| \|h_1\|\right]^{2l-\varepsilon} + c_3 \left[|\beta| \|h_2\|\right]^{2l-\varepsilon}\right),$$

where constant $c_1 > 0$, $c_2 > 0$, $c_3 > 0$, $0 < \varepsilon < 1$ depend only of the function f.

Lemma 1. It exists a sequence of positive numbers $\alpha_n \leq \lambda_n$ such that we have

$$\int_{H} \left[\prod_{n=1}^{\infty} \left(1 + \alpha_n t \frac{\|x\|^2}{p(Sx)^2} \right) \right]^2 \exp\left(-\frac{\|x\|^4}{p(Sx)^2} \right) \mu(\mathrm{d}x) \le \mathrm{e}^t$$

for all real $t \ge 0$ and for all self-adjoint strictly positive bounded operator $S: H \to H$.

Proof. Let us take $a_0 = 1$,

$$a_n = n^{-2n} \left[1 + \sum_{k=1}^n \int_H \frac{1}{\|x\|^{2k}} \,\mu(\mathrm{d}x) \right]^{-1},$$

 $\alpha_n = 2^{-2n} a_{2n}$ for all natural number *n*. We have

$$\left[\prod_{n=1}^{\infty} \left(1 + \alpha_n t \frac{\|x\|^2}{p(Sx)^2}\right)\right]^2 \exp\left(-\frac{\|x\|^4}{p(Sx)^2}\right) \le \\ \le \sum_{n=0}^{\infty} a_n \left(t \frac{\|x\|^2}{p(Sx)^2}\right)^n \exp\left(-\frac{\|x\|^4}{p(Sx)^2}\right) \le \\ \le \sum_{n=0}^{\infty} a_n t^n \frac{1}{\|x\|^{2n}} \left[\frac{\|x\|^4}{p(Sx)^2}\right)^n \exp\left(-\frac{\|x\|^4}{p(Sx)^2}\right) \le \sum_{n=0}^{\infty} n^n a_n t^n \frac{1}{\|x\|^{2n}} \cdot \frac{1}{\|x\|^{2n}} \left[\frac{\|x\|^4}{p(Sx)^2}\right]^n \exp\left(-\frac{\|x\|^4}{p(Sx)^2}\right) \le \sum_{n=0}^{\infty} n^n a_n t^n \frac{1}{\|x\|^{2n}} \cdot \frac{1}{\|x\|^{2$$

Hence

$$\int_{H} \left[\prod_{n=1}^{\infty} \left(1 + \alpha_{n} t \frac{\|x\|^{2}}{p(Sx)^{2}} \right) \right]^{2} \exp\left(-\frac{\|x\|^{4}}{p(Sx)^{2}} \right) \mu(\mathrm{d}x) \leq \\ \leq \sum_{n=0}^{\infty} n^{n} a_{n} t^{n} \int_{H} \frac{1}{\|x\|^{2n}} \mu(\mathrm{d}x) \leq \sum_{n=0}^{\infty} n^{-n} t^{n} \leq \mathrm{e}^{t} \,.$$

Lemma 2. Let d be a natural number. Then it exists a self-adjoint strictly positive bounded operator S such that

$$Se_n = \gamma_n e_n, \qquad \alpha_n \ge \gamma_n \ge \gamma_{n+1} > 0,$$

 $\frac{\|P_n x\|^2}{(p(SP_n x))^2} \le 2 \frac{\|x\|^2}{(p(Sx))^2}$

 $\mathbf{2}$

and

$$\int_{H} \exp\left(\left[\frac{\|T^{-1}Sx\|}{p(Sx)}\right]^{d}\right) \mu(\mathrm{d}x) \le +\infty.$$

Proof. Let $\gamma_1 = \alpha_1$, and $I = \exp\left(\left[\frac{\lambda_1^{-1}}{p(e_1)}\right]^d\right)$. It is easy to see that

$$I = \frac{1}{\sqrt{2\pi\lambda_1}} \int_{-\infty}^{+\infty} \exp\left(\left[\frac{\lambda_1^{-1}\gamma_1|x_1|}{p(\gamma_1 x_1 e_1)}\right]^d\right) \exp\left(-\frac{x_1^2}{2\lambda_1}\right) \mathrm{d}x_1 \,.$$

Suppose that we have $\gamma_1, \ldots, \gamma_n$ such that $\alpha_1 \ge \gamma_1 > 0, \ldots, \alpha_n \ge \gamma_n > 0$, $\gamma_1 \ge \gamma_2 \ge \ldots \ge \gamma_n > 0$ and

$$\begin{split} I_k &= \frac{1}{\sqrt{(2\pi)^n \lambda_1 \dots \lambda_{k-1} \lambda_k}} \times \\ &\times \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \left| \exp\left(\left[\frac{\sqrt{(\lambda_1^{-1} \gamma_1 x_1)^2 + \dots + (\lambda_k^{-1} \gamma_k x_k)^2}}{p(\gamma_1 x_1 e_1 + \dots + \gamma_k x_k e_k)} \right]^d \right) - \\ &- \exp\left(\left[\frac{\sqrt{(\lambda_1^{-1} \gamma_1 x_1)^2 + \dots + (\lambda_{k-1}^{-1} \gamma_{k-1} x_{k-1})^2}}{p(\gamma_1 x_1 e_1 + \dots + \gamma_{k-1} x_{k-1} e_{k-1})} \right]^d \right) \right| \times \\ &\times \exp\left(-\frac{x_1^2}{2\lambda_1} - \dots - \frac{x_{k-1}^2}{2\lambda_{k-1}} - \frac{x_k^2}{2\lambda_k} \right) \mathrm{d}x_1 \dots \mathrm{d}x_{k-1} \mathrm{d}x_k < 2^{-k} \end{split}$$

for all k = 2, 3, ..., n.

Let us find γ_{n+1} . Denoted by

$$I_{n+1}(t) = \frac{1}{\sqrt{(2\pi)^{n+1}\lambda_1 \dots \lambda_n \lambda_{n+1}}} \times \\ \times \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \left| \exp\left(\left[\frac{\sqrt{(\lambda_1^{-1} \gamma_1 x_1)^2 + \dots + (\lambda_{n+1}^{-1} t x_{n+1})^2}}{p(\gamma_1 x_1 e_1 + \dots + \gamma_n x_n e_n + t x_{n+1} e_{n+1})} \right]^d \right) - \\ - \exp\left(\left[\frac{\sqrt{(\lambda_1^{-1} \gamma_1 x_1)^2 + \dots + (\lambda_n^{-1} \gamma_n x_n)^2}}{p(\gamma_1 x_1 e_1 + \dots + \gamma_n x_n e_n)} \right]^d \right) \right| \times \\ \times \exp\left(-\frac{x_1^2}{2\lambda_1} - \dots - \frac{x_n^2}{2\lambda_n} - \frac{x_{n+1}^2}{2\lambda_{n+1}} \right) dx_1 \dots dx_n dx_{n+1}$$

for all $t \geq 0$.

By Lebesgue's dominated convergence theorem we get

$$\lim_{t \to +0} I_{n+1}(t) = 0 \, .$$

E.T. Shavgulidze

Thus it exist t > 0 such that $I_{n+1}(t) < 2^{-(n+1)}, t \le \min(a_{n+1}, \gamma_n)$ and

$$\frac{\|P_n x\|^2}{(p(SP_n x))^2} \le \frac{2 - \frac{1}{n+1}}{2 - \frac{1}{n}} \frac{\|P_{n+1} x\|^2}{(p(SP_{n+1} x))^2}.$$

We can take $\gamma_{n+1} = t$. Thus we find $\gamma_1, \ldots, \gamma_n, \ldots$.

Let us take d = 6. By lemma 2 we find self-adjoint strictly positive bounded operator S such that $Se_n = \gamma_n e_n$. Let $t = t_1 + t_2 i \in \mathbb{C}$, $t_1, t_2 \in \mathbb{R}$, $t_1 \ge 0$, $t \ne 0$ and $x = \sum_{n=1}^{\infty} x_n e_n \in H$, $x \ne 0$. We suppose

$$a(x) = \left(1 + t_1 \frac{\|x\|^2}{p(Sx)^2}\right)^2 + \left(t_2 \frac{\|x\|^2}{p(Sx)^2}\right)^2,$$

$$A_1(x) = \sum_{n=1}^{\infty} \frac{1 + t_1 \frac{\|x\|^2}{p(Sx)^2}}{a(x)} \gamma_n x_n e_n$$

and

$$A_2(x) = \sum_{n=1}^{\infty} \frac{t_2 \frac{\|x\|^2}{p(Sx)^2}}{a(x)} \gamma_n x_n e_n \,.$$

We take

$$J_1(t,x) = \prod_{n=1}^{\infty} \left(1 + \gamma_n t \frac{\|x\|^2}{p(Sx)^2} \right)$$

and

$$J_{2}(t,x) = 1 + 2t \frac{(x,A_{1}(x)) + i(x,A_{2}(x))}{p(Sx)^{2}} - \frac{\|x\|^{2}}{p(Sx)^{2l+2}} \Big[Q(SA_{1}(x),Sx,\dots,Sx) + iQ(SA_{2}(x),Sx,\dots,Sx) \Big].$$

Theorem 1. If a function f belongs to the space F(H), then for all $\nu, \nu \in \mathbf{C}$, $\nu \neq 0$, upsilon $\neq 0$, Re $\nu \geq 0$, Re $\nu \geq 0$ the Feynman integral

$$\int_{H} f(x) \exp\left(-\nu(p(x))^{2l}\right) \exp\left(-\frac{\nu}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x$$

exists and it can also be represented by

$$\begin{split} \int_{H} f(x) \exp\left(-\nu(p(x))^{2l}\right) \exp\left(-\frac{\upsilon}{2} \left(T^{-1}x, x\right)\right) \mathrm{d}x &= \int_{H} f\left(\sigma x + \omega \,\frac{\|x\|^{2}}{(p(Sx))^{2}} Sx\right) \times \\ & \times \exp\left(-\nu\left(p(\sigma x + \omega \,\frac{\|x\|^{2}}{(p(Sx))^{2}} Sx\right)\right)^{2l}\right) \exp\left(-\frac{\omega}{\sigma} \,\frac{\|x\|^{2}}{(p(Sx))^{2}} (x, T^{-1}Sx)\right) \times \\ & \times \exp\left(-\frac{\upsilon\omega^{2}}{2} \,\frac{\|x\|^{4}}{(p(Sx))^{4}} \left(Sx, T^{-1}Sx\right)\right) J_{1}\left(\frac{\omega}{\sigma}, x\right) J_{2}\left(\frac{\omega}{\sigma}, x\right) \mu(\mathrm{d}x), \end{split}$$

where
$$\sigma = v^{-1/2}$$
, $\omega = 2^{4l} r v^{-1/(2l)}$, $r \in \mathbf{R}$, $r > |\nu| + |\nu| + 1$.
Proof. Let $\nu, v \in \mathbf{R}$, $\nu > 0$, $v > 0$, $r \in \mathbf{R}$, $r > |\nu| + |v| + 1$ and $\sigma = v^{-1/2}$,
 $\omega = 2^{4l} r v^{-1/(2l)}$.
Let $x' = \sigma x + \omega \frac{||x||^2}{(p(Sx))^2} Sx$ for all $x \in H_n$. We have

$$\frac{1}{\sqrt{(\frac{2\pi}{v})^n \lambda_1 \dots \lambda_n}} \int_{H_n} f(x') \exp\left(-\nu(p(x'))^{2l}\right) \exp\left(-\frac{v}{2}(T^{-1}x', x')\right) dx' =$$

$$= \frac{1}{\sqrt{(2\pi)^n \lambda_1 \dots \lambda_n}} \int_{H_n} f\left(\sigma x + \omega \frac{||x||^2}{(p(Sx))^2} Sx\right) \exp\left(-\nu\left(p(\sigma x + \omega \frac{||x||^2}{(p(Sx))^2} Sx)\right)^{2l}\right)$$

$$\times \exp\left(-\frac{\omega}{\sigma} \frac{||x||^2}{(p(Sx))^2}(x, T^{-1}Sx) - \frac{v\omega^2}{2} \frac{||x||^4}{(p(Sx))^4}(Sx, T^{-1}Sx)\right) \times$$

$$\times J_1\left(\frac{\omega}{\sigma}, x\right) J_2\left(\frac{\omega}{\sigma}, x\right) e^{-(T^{-1}x, x)/2} dx.$$

Taking the limit for $n \to \infty$ we prove the theorem 1. The convergence follow from the lemma 1 and the lemma 2.

We say that a function

$$v: [0, +\infty) \times \mathbf{R}^n \to \mathbf{R}$$

satisfies condition (A_k) if

$$v(t,q) = v_1(t,q) + v_2(t,q),$$

where for all $t \ge 0$ the function $v_1(t, .)$ is nondegenerate homogeneous polynomial of degree 2k for argument q and the function $v_2(t, \cdot)$ is polynomial of degree less then 2k.

Consider Schrödinger equation

$$\frac{\partial u(t,q)}{\partial t} = \frac{1}{2} i\Delta u(t,q) - iv(t,q)u(t,q)$$
(1)

with a potential

$$v: [0, +\infty) \times \mathbf{R}^n \to \mathbf{R}.$$

Denote by E_t the Banach space of all continuous mapping $x : [0, t] \in \mathbf{R}^n$ such that x(0) = x(t) = 0.

Theorem 2. If the potential v is continuous and satisfy condition (A_k) for a naturel k then it exists for the Schrödinger equation (1) the Green's function

$$U(t, q_0, q) = \int_{E_t} \exp\left[-i \int_0^t v\left(\tau, x(\tau) + q_0 + \frac{\tau}{t} (q - q_0)\right) d\tau\right] \times \\ \times \exp\left[-\frac{i \|q - q_0\|_{\mathbf{R}^n}^2}{2t} + \frac{i}{t} \int_0^t (q - q_0, x(\tau))_{\mathbf{R}^n} d\tau\right] \exp\left[-\frac{i}{2} \int_0^t \|x'(\tau)\|_{\mathbf{R}^n}^2 d\tau\right] dx$$

E.T. Shavgulidze: Feynman path integrals for exponents of polynomially growing functionals

and

$$U\left(0,q_{0},q\right)=\delta(q-q_{0}).$$

Remark. The Green's function $U(t, q_0, q)$ defines special self-adjoint extension for the Schrödinger differential operator $\triangle + v$ in the case that this operator is not essentially self-adjoint.

Consider on the Hilbert space Q = H the Cauchy problem

$$\frac{\partial u(t,q)}{\partial t} = \frac{\mathrm{i}}{2} \Delta_T u(t,q) - \mathrm{i} V(t,q) u(t,q)$$
(2)

of Schrödinger equation with a potential $V:[0,+\infty)\times Q\to {\bf R}$ and with a initial condition

$$u(0,q) = u_0(q-q_0) \quad (\forall q \in Q) ,$$
 (3)

where $u_0: Q \to \mathbf{C}$,

$$\Delta_T = \sum_{n=1}^{\infty} \lambda_n \frac{\partial^2}{\partial q_n^2} \qquad q = \sum_{n=1}^{\infty} q_n e_n \in Q \,.$$

Denote by F_t the Banach space of all continuous mapping $\xi : [0, t] \in Q$ such that $\xi(t) = 0$.

Let $V(t,q) = V_1(t,q,...,q) + V_2(t,q)$, where

$$V_1: [0, +\infty) \times Q \times \cdots \times Q \in \mathbf{R},$$

 $V_2: [0, +\infty) \times Q \to \mathbf{R}$ are continuous functions and for all $t \ge 0$ the function $V_1(t, ..., ..)$ is 2k-linear symmetric nondegenerate and the function $V_1(t, ..)$ is polynomial of degree less then 2k for argument q.

Theorem 3. If a function u_0 belongs to the space F(Q) then it exists a solution of the Cauchy problem of the Schrödinger equation (2), (3) and

$$\begin{split} u(t,q) &= \int_{F_t} u_0 \big(\xi(0) + q \big) \exp \bigg(-\mathrm{i} \int_0^t V \big(\tau, \xi(\tau) + q \big) \mathrm{d} \tau \bigg) \times \\ &\times \exp \bigg(-\mathrm{i} \int_0^t \Big\| T^{-1/2} \xi'(\tau) \Big\|_Q^2 \, \mathrm{d} \tau \bigg) \mathrm{d} \xi \,. \end{split}$$

References

- O.G. Smolyanov and E.T. Shavgulidze: *Path integrals*. Moskov. Gos. Univ., Moscow, 1990.
- [2] E.T. Shavgulidze: Russ. Math. Dokl. 348 (1996) 743.